

Content



0. Introduction

1. Differential equations

1.1 Classical equation of motion (classical mechanics, pendulum)

1.2 Partial differential equation relaxation methods (electromagnetism, diffusion)

1.3 Partial differential equation in space-time (traffic flow, tsunami)

2. Eigenvalue problem

2.1 Schrödinger equation and Hamiltonian (Harmonic oscillator, wave package)

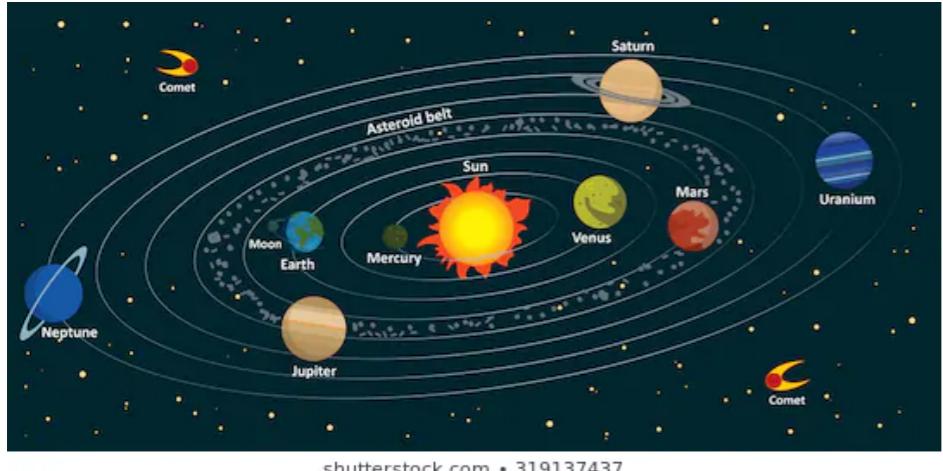
2.2 Quantum lattice model and Hibert space (Heisenberg model)

2.3 Exact diagonalization of spin chain (Spin wave, Haldane conjecture, topology)

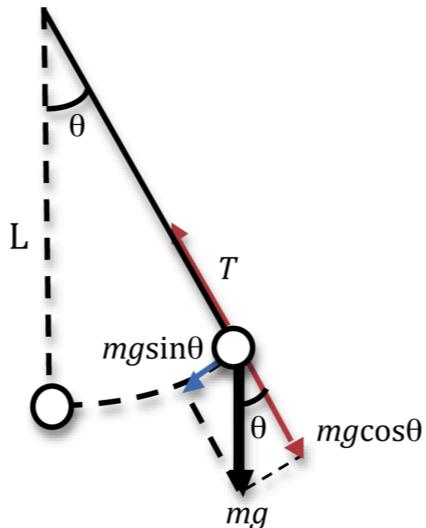
2.4 Matrix product state and density matrix renormalization group (DMRG)

Differential equations

- Initial value problems: time-dependent equations with given initial conditions



Solar system



Pendulum



浮世绘，葛饰北斋，神奈川冲浪里

- Boundary value problems: differential equations with specific boundary values



Eigenvalue problems

CAN ONE HEAR THE SHAPE OF A DRUM?

MARK KAC, The Rockefeller University, New York

To George Eugene Uhlenbeck on the occasion of his sixty-fifth birthday

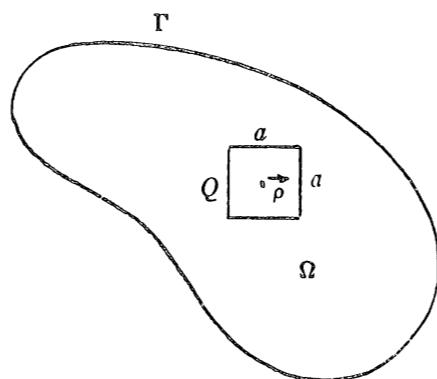
Am. Math. Mon. 73, 1 (1966)

Eigenvalues of Dirichlet problem for Laplacian

$$\frac{1}{2} \nabla^2 U + \lambda U = 0 \text{ in } \Omega,$$

$$U = 0 \text{ on } \Gamma.$$

Length of circumference



$$\sum_{n=1}^{\infty} e^{-\lambda_n t} \sim \frac{|\Omega|}{2\pi t} - \frac{L}{4} \frac{1}{\sqrt{2\pi t}} + \frac{1}{6} \frac{(1-r)}{(1-r)\frac{1}{6}}$$

Number of holes

$$(1-r)\frac{1}{6}$$

Classical equation of motion

Differential equations

$$\dot{\vec{v}}_i(t) = \vec{a}_i(\vec{x}_0(t), \dots, \vec{x}_{N-1}(t), \vec{v}_0(t), \dots, \vec{v}_{N-1}(t), t)$$

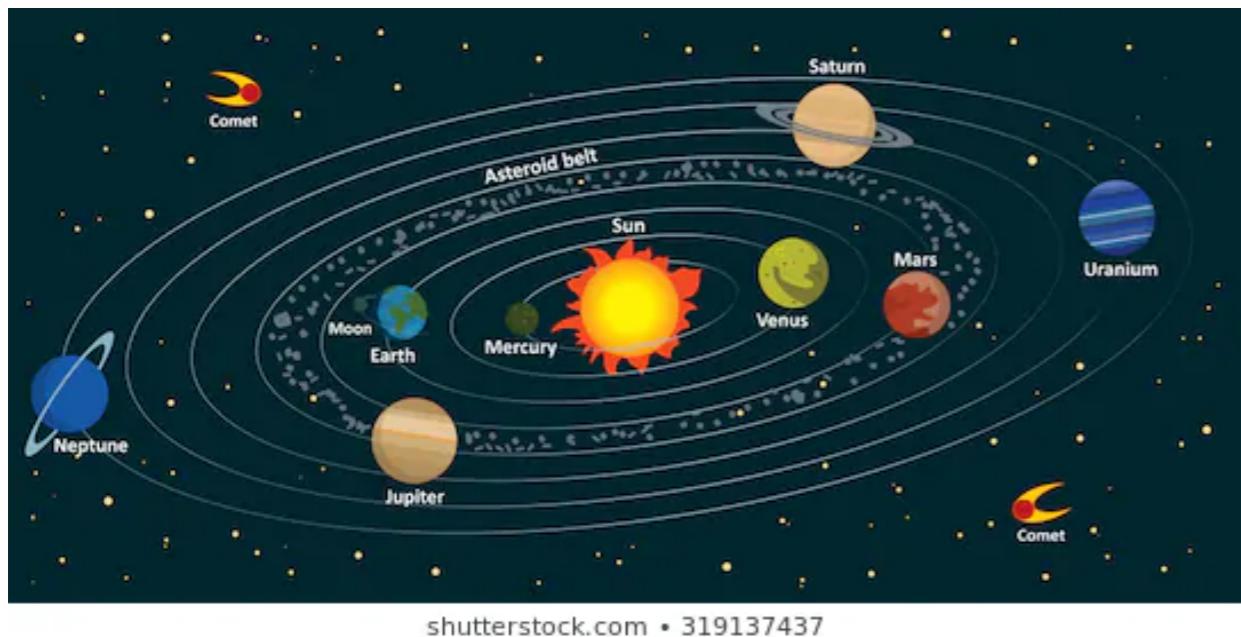
State of dynamical system

$$\dot{\vec{x}}_i(t) = \vec{v}_i(t) \quad i = 0, 1, \dots, N-1$$

Gravitation (such as solar system)

$$\vec{a}_i(\vec{x}_0(t), \dots, \vec{x}_{N-1}(t)) = G \sum_{j \neq i} \frac{m_j}{|\vec{x}_j(t) - \vec{x}_i(t)|^3} [\vec{x}_j(t) - \vec{x}_i(t)]$$

G Gravitational constant



Discretization

$$t = t_0, t_1, t_2, \dots \quad \tau = t_{n+1} - t_n$$

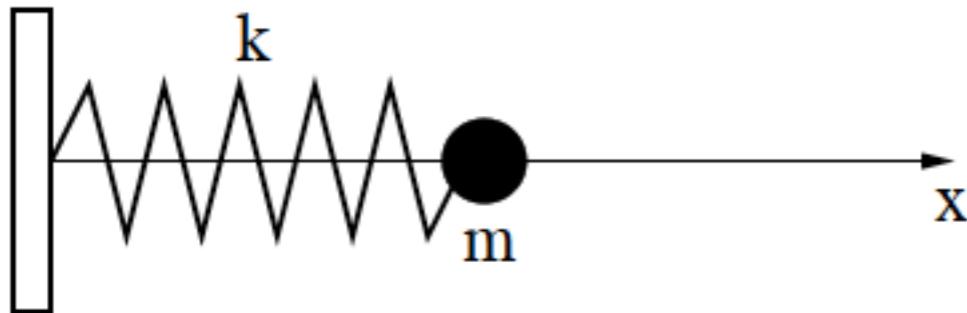
$$\vec{x}(t) = \begin{pmatrix} \vec{x}_0(t) \\ \vec{x}_1(t) \\ \vdots \\ \vec{x}_{N-1}(t) \end{pmatrix} \quad \vec{v}(t) = \begin{pmatrix} \vec{v}_0(t) \\ \vec{v}_1(t) \\ \vdots \\ \vec{v}_{N-1}(t) \end{pmatrix}$$

$$\dot{\vec{v}}(t) = \vec{a}(\vec{x}(t), \vec{v}(t), t)$$

$$\dot{\vec{x}}(t) = \vec{v}(t)$$

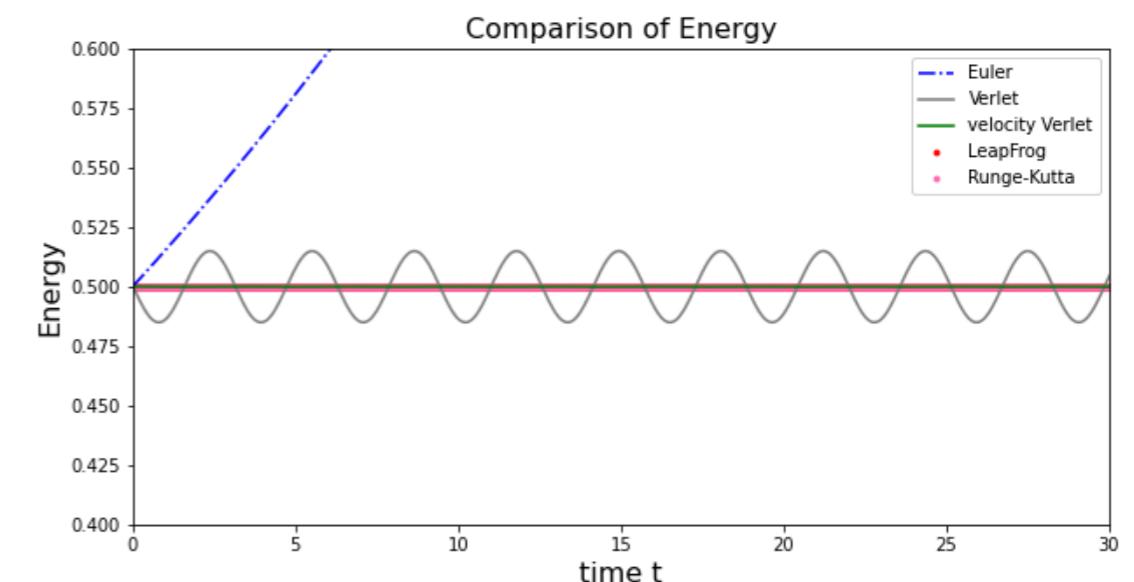
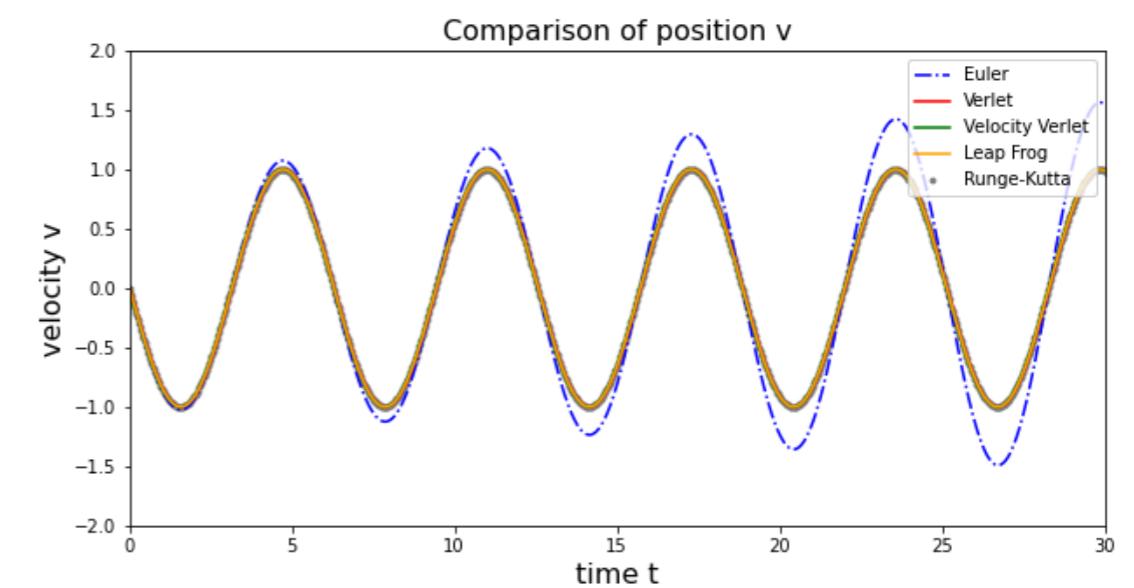
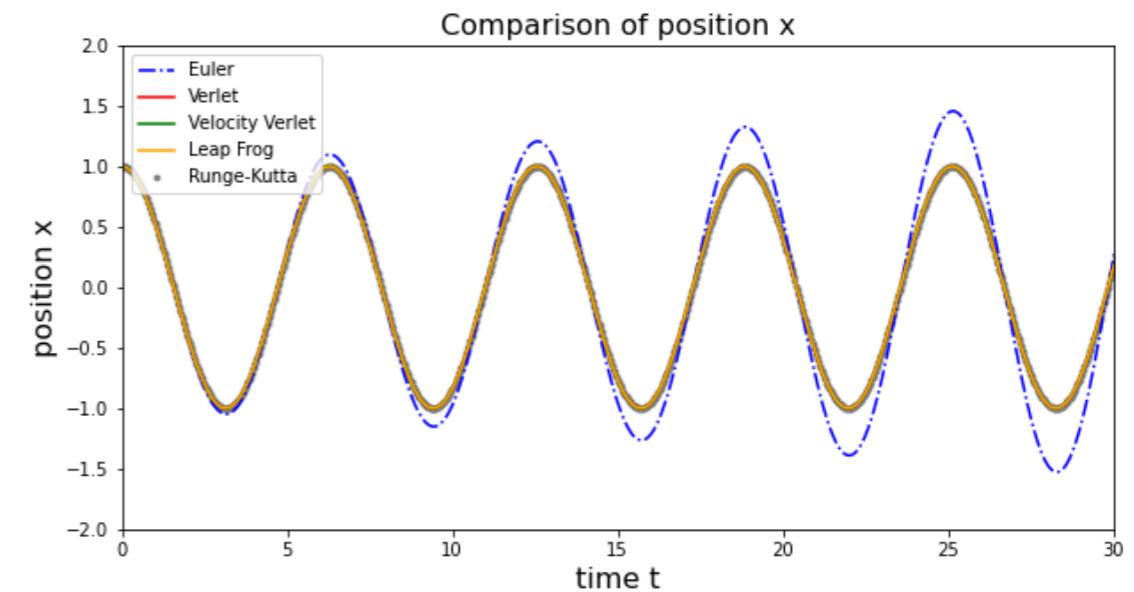
Classical equation of motion

Harmonic Oscillator



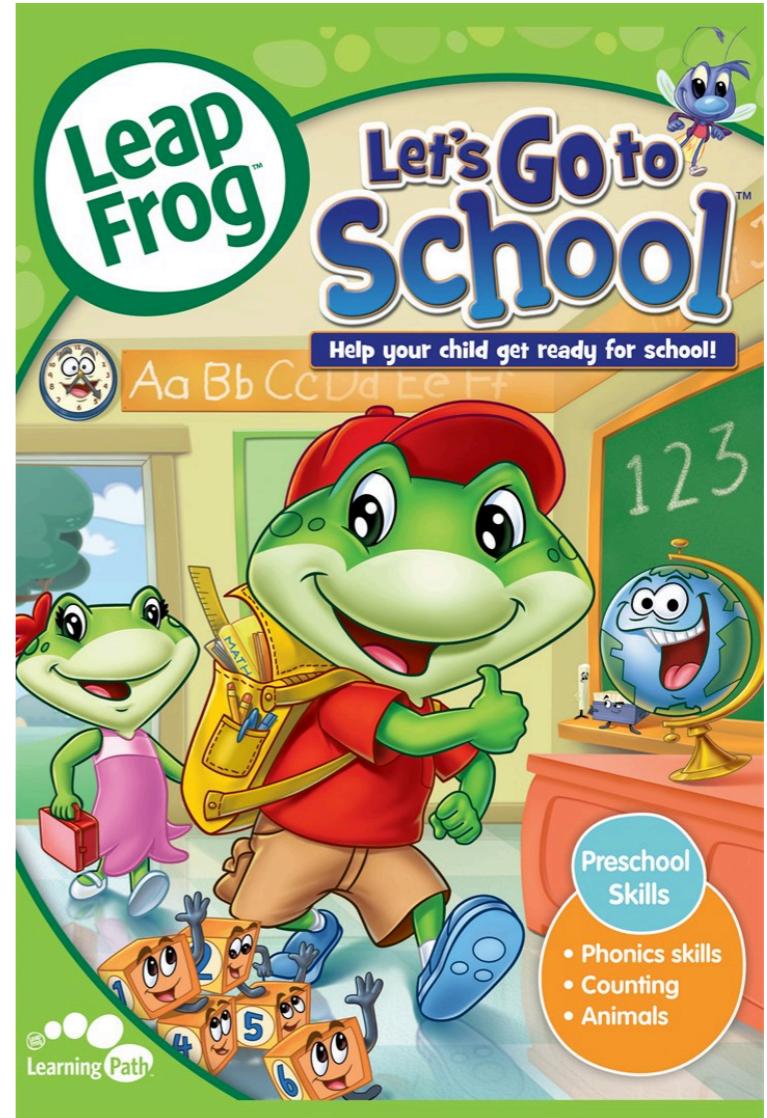
$$\vec{a}(t) = \dot{\vec{v}}(t) = -\frac{k}{m}x(t)$$

$$\dot{\vec{x}}(t) = \vec{v}(t)$$



Initial value problem

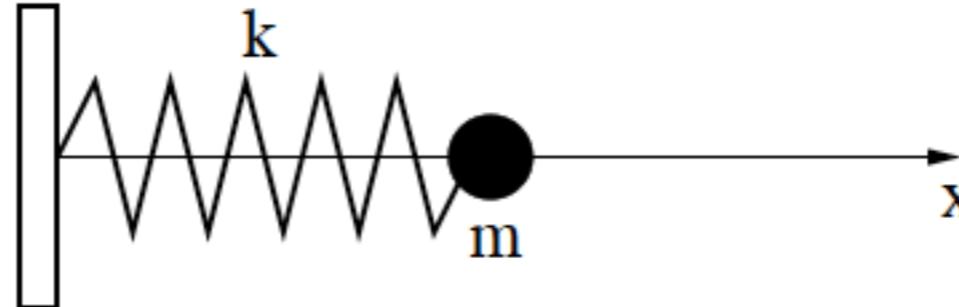
- Euler (τ^2)
- Verlet (τ^4 , time-reversal)
- Velocity Verlet (variant of Verlet 1)
- Leap-frog (variant of Verlet 2, less roundoff error)
- Runge-Kutta (τ^5 , workhorse)



LeapFrog Enterprises, Inc. is the leader in innovative solutions that encourage a child's curiosity and love of learning throughout their early developmental journey.

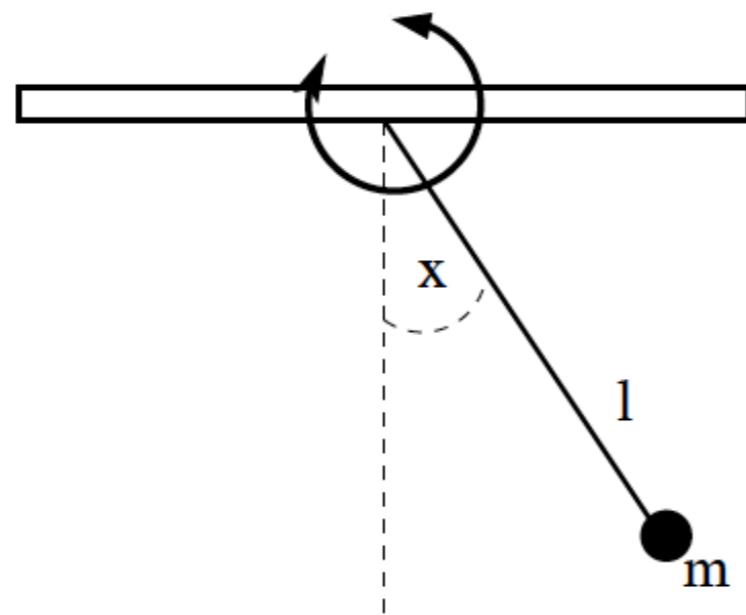
Classical equation of motion

Harmonic Oscillator



$$\dot{v}(t) = -\frac{k}{m}x$$
$$\dot{x}(t) = v$$

Driven Pendulum



$$\gamma = 0 \quad \Omega = 0$$

goes back to harmonic oscillator at small x

$$g = 9.81 \text{ m/s}^2$$

Soliton solution

γ friction coefficient Stoke's friction

Q strength of periodic driving force

Ω driving frequency

$$\dot{v}(t) = -\frac{g}{l} \sin(x) - \gamma \dot{x} + Q \sin(\Omega t)$$

$$\dot{x}(t) = v$$

Many interesting videos

https://www.youtube.com/watch?v=_TSp6KkMbP4

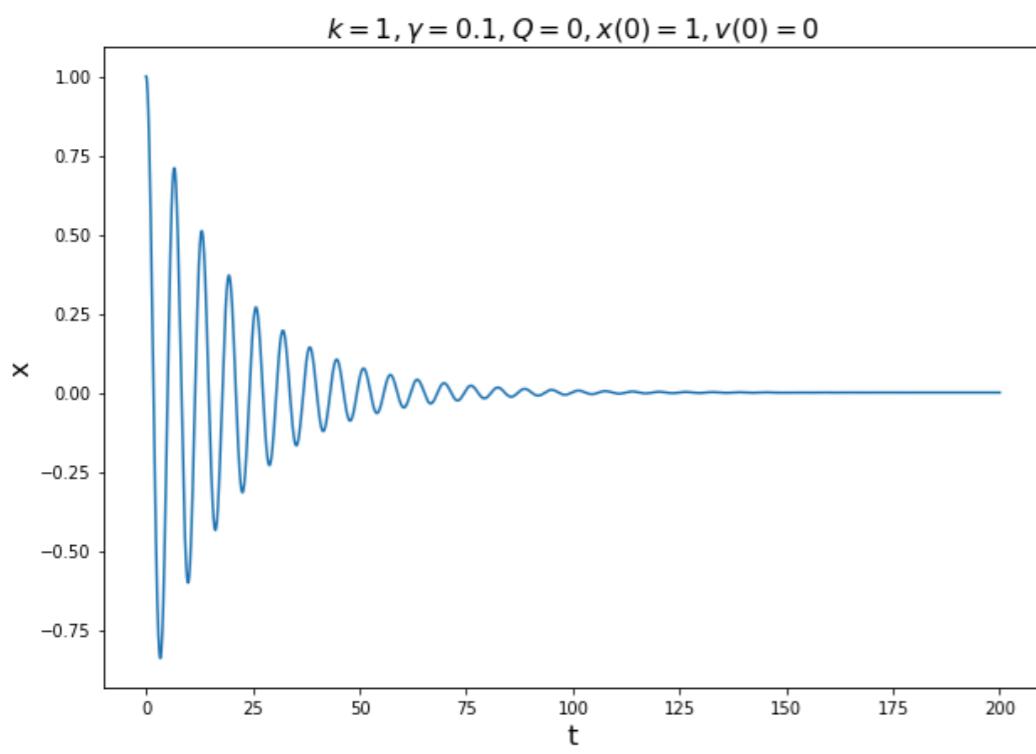
$$\ddot{x} = -k \sin(x) - \gamma \dot{x} + Q \sin(\Omega t) \quad k = g/l = 1$$

$\gamma \neq 0$ and $Q \neq 0$ no analytical solution

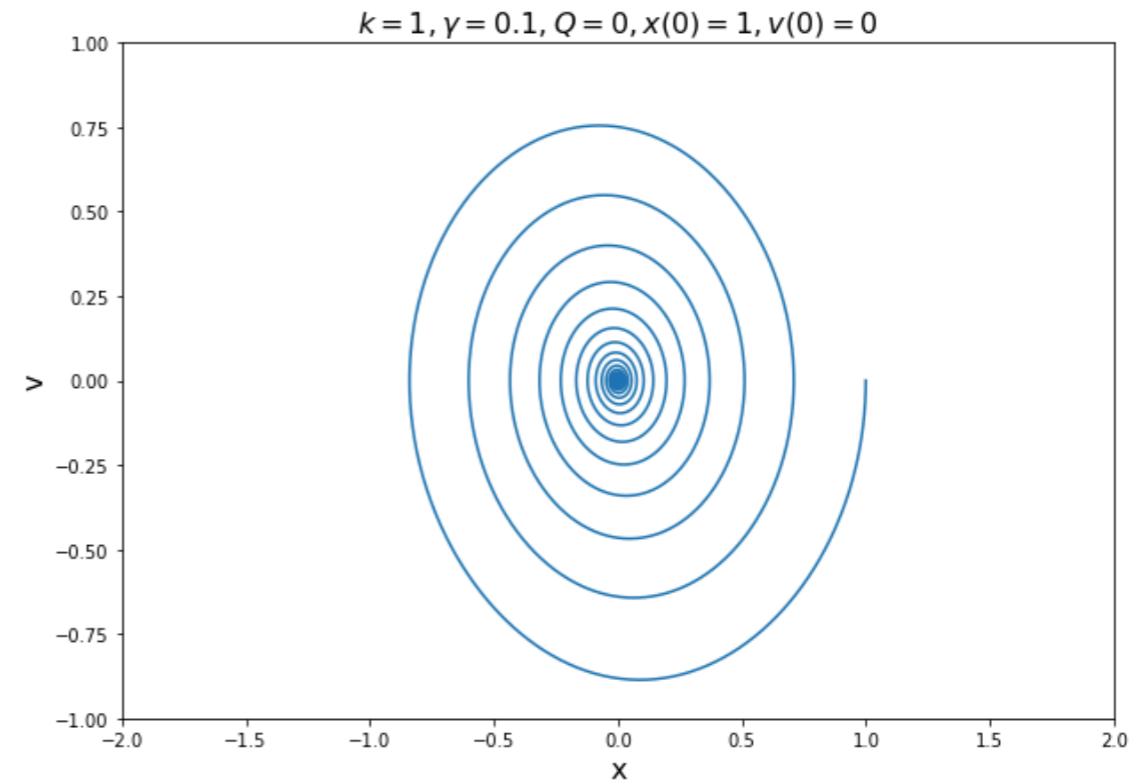
$$x \in [-\pi, \pi]$$

$$\gamma \neq 0 \quad \text{and} \quad Q = 0$$

With friction only, damping



In phase space, attractor



$$\gamma \neq 0 \quad \text{and} \quad Q \neq 0$$

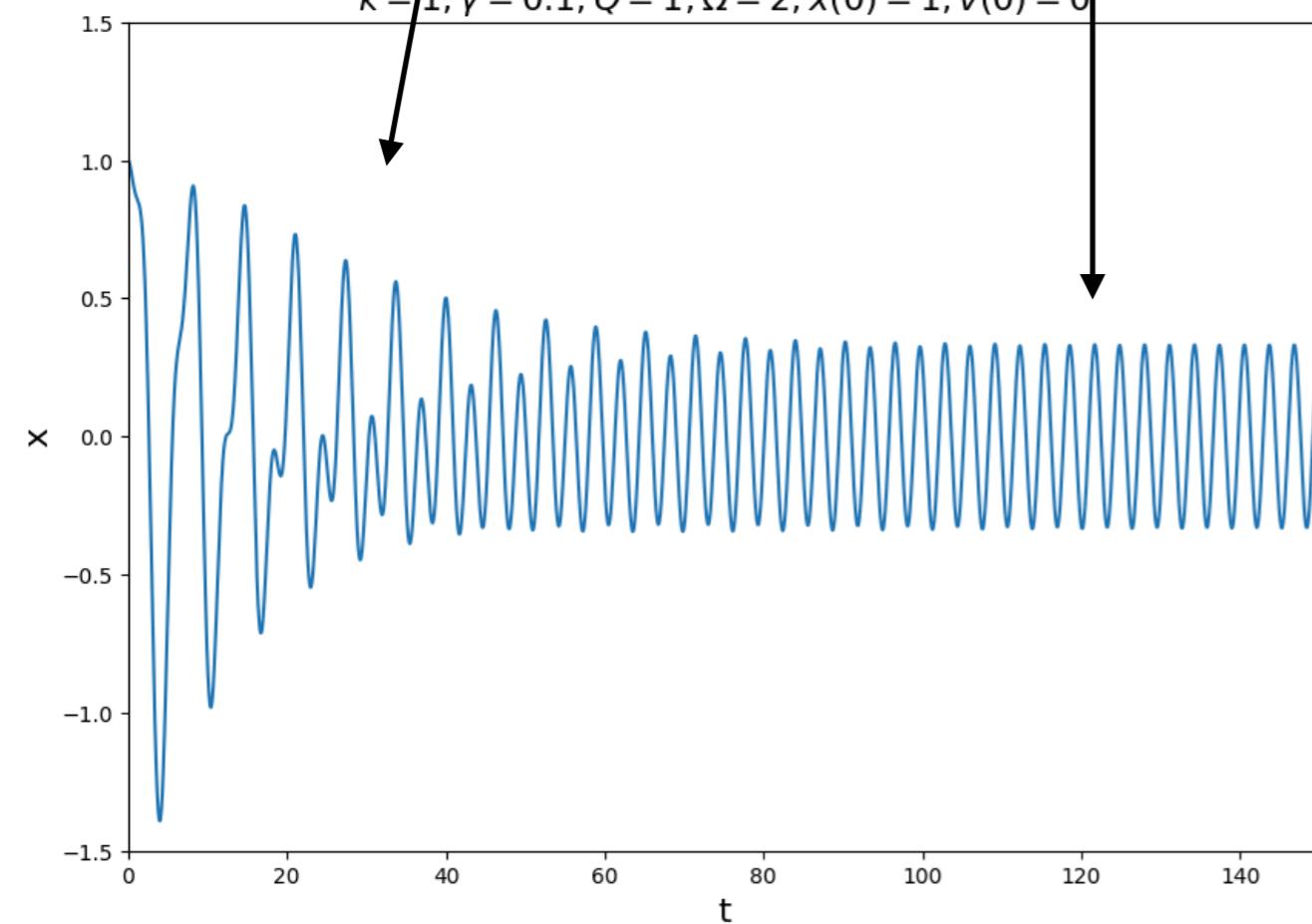
The original period

$$T = 2\pi \sqrt{\frac{m}{k}} = 2\pi$$

Initial damping and transient behaviour

Driving force with frequency $\Omega = 2$

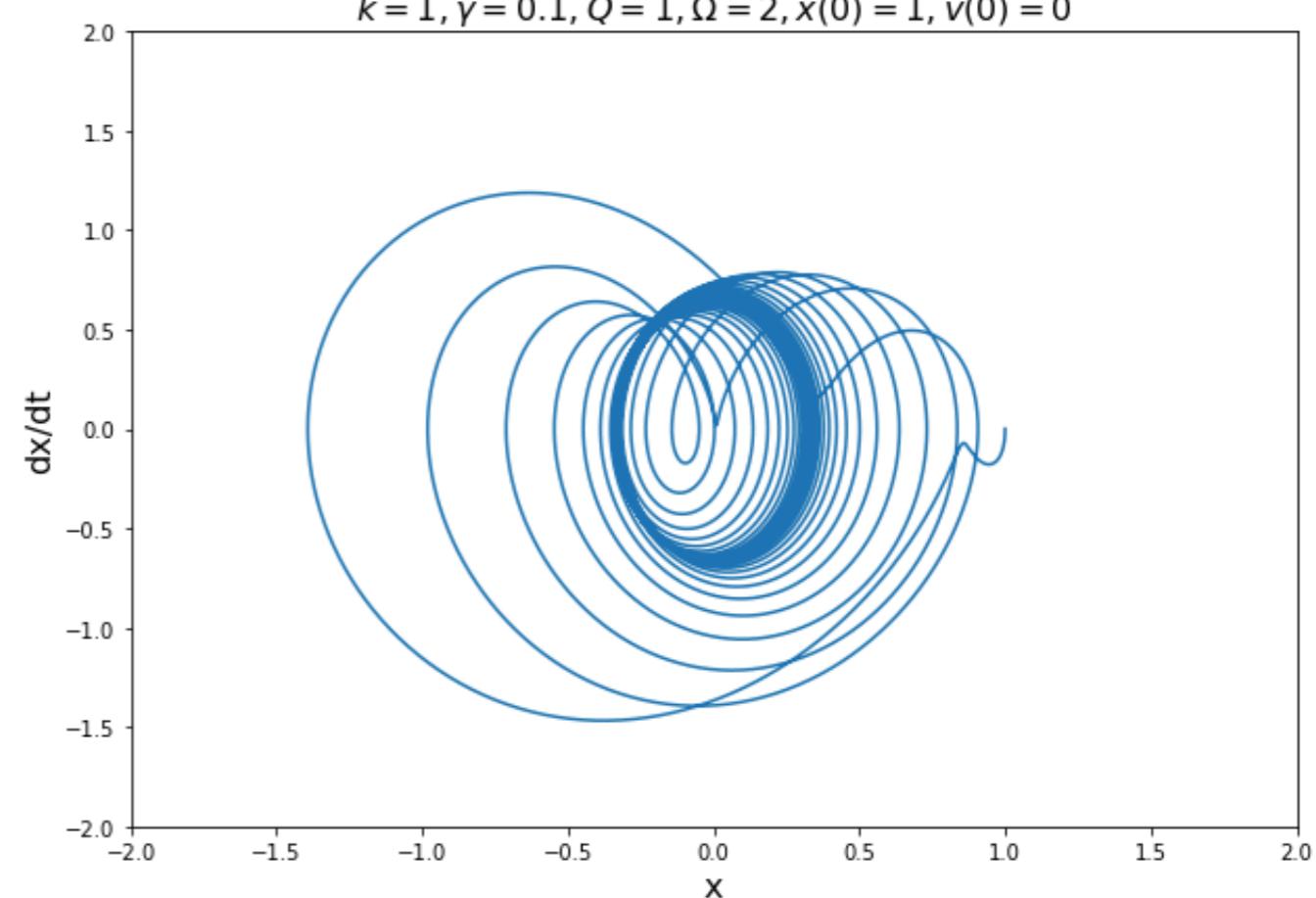
$$k=1, \gamma=0.1, Q=1, \Omega=2, x(0)=1, v(0)=0$$



In phase space, cyclic attractor

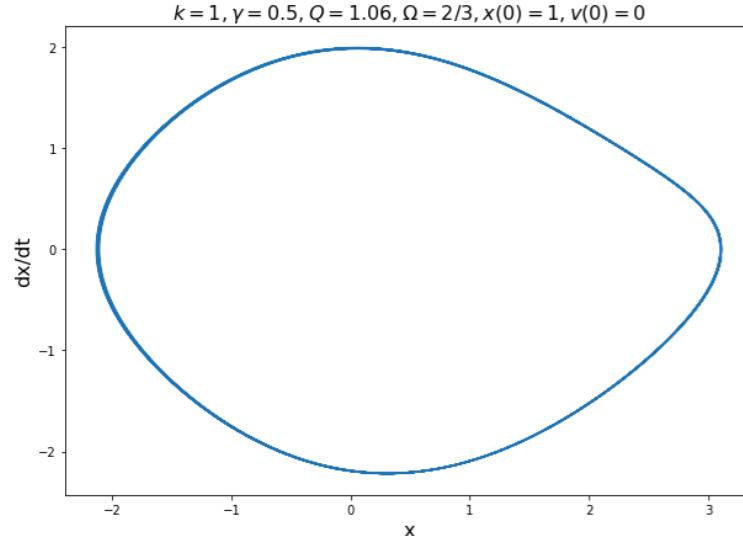
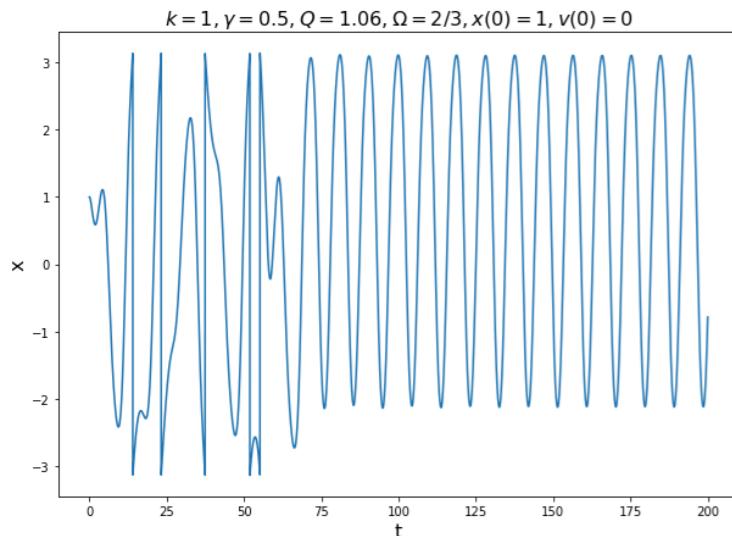
$$\text{Period} \quad T_\Omega = \frac{2\pi}{\Omega}$$

$$k=1, \gamma=0.1, Q=1, \Omega=2, x(0)=1, v(0)=0$$



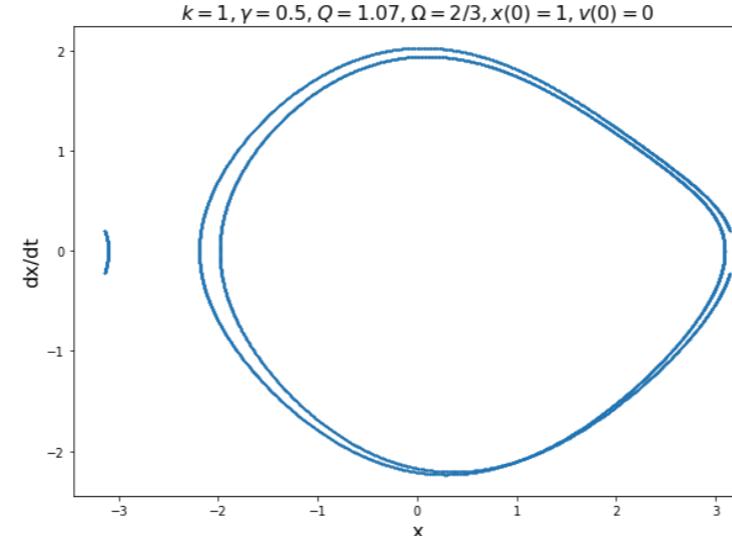
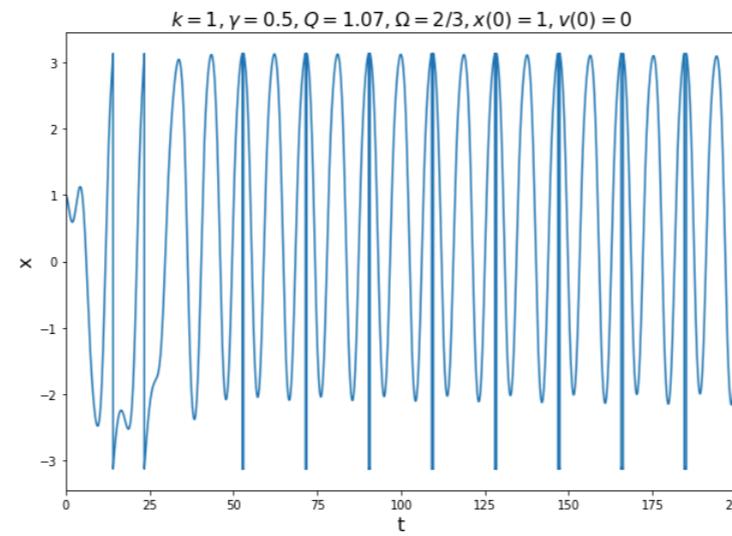
Asymmetric attractor
Spontaneous breaking of reflection symmetry

$$(x, v) \neq (-x, -v) \quad Q = 1.06, \quad T = T_\Omega$$



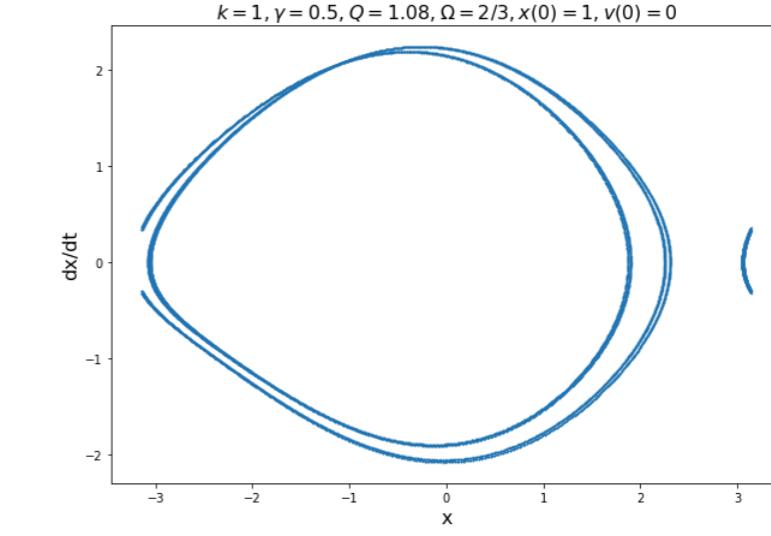
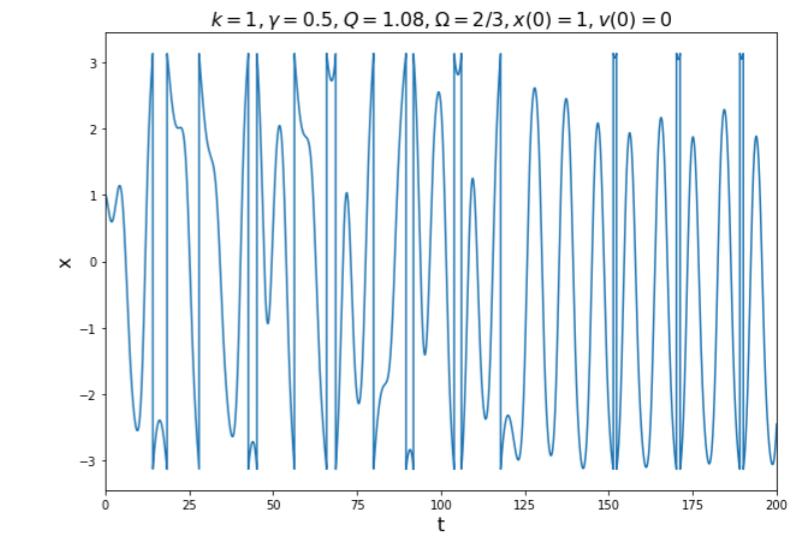
Period doubling

$$Q = 1.07, \quad T = 2T_\Omega$$



Further period doubling

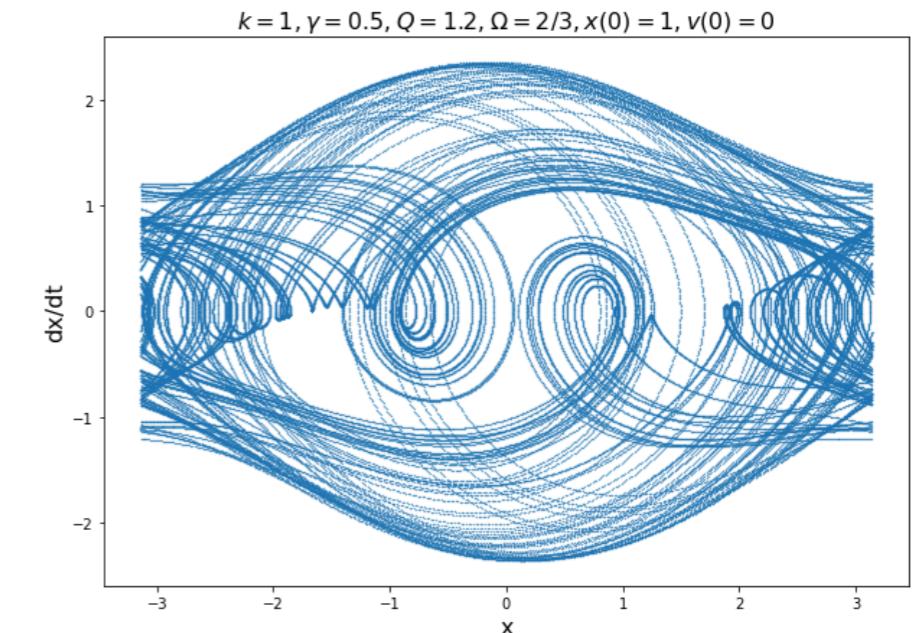
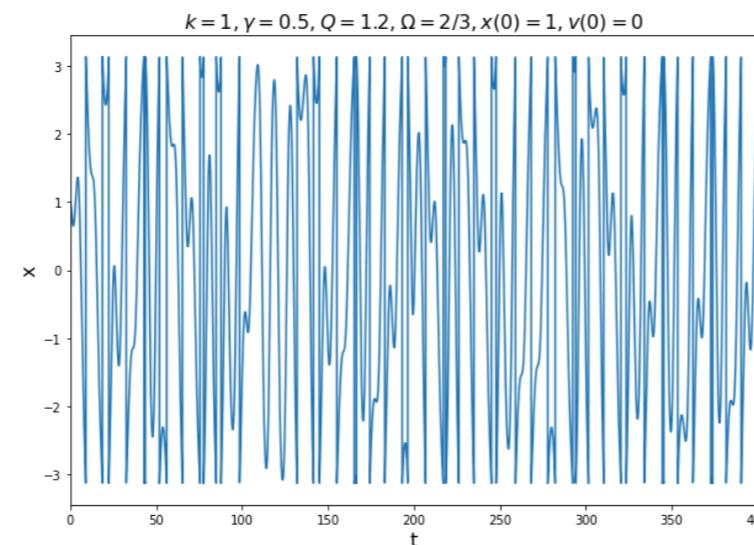
$$Q = 1.08, \quad T = 4T_\Omega$$



Chaotic attractor

$$Q = 1.2$$

The motion never closes,
the attractor fills a finite
region of the phase space



Chaos theory

https://en.wikipedia.org/wiki/Chaos_theory

https://en.wikipedia.org/wiki/Chaos_theory#/media/File:Double-compound-pendulum.gif

Butterfly effect

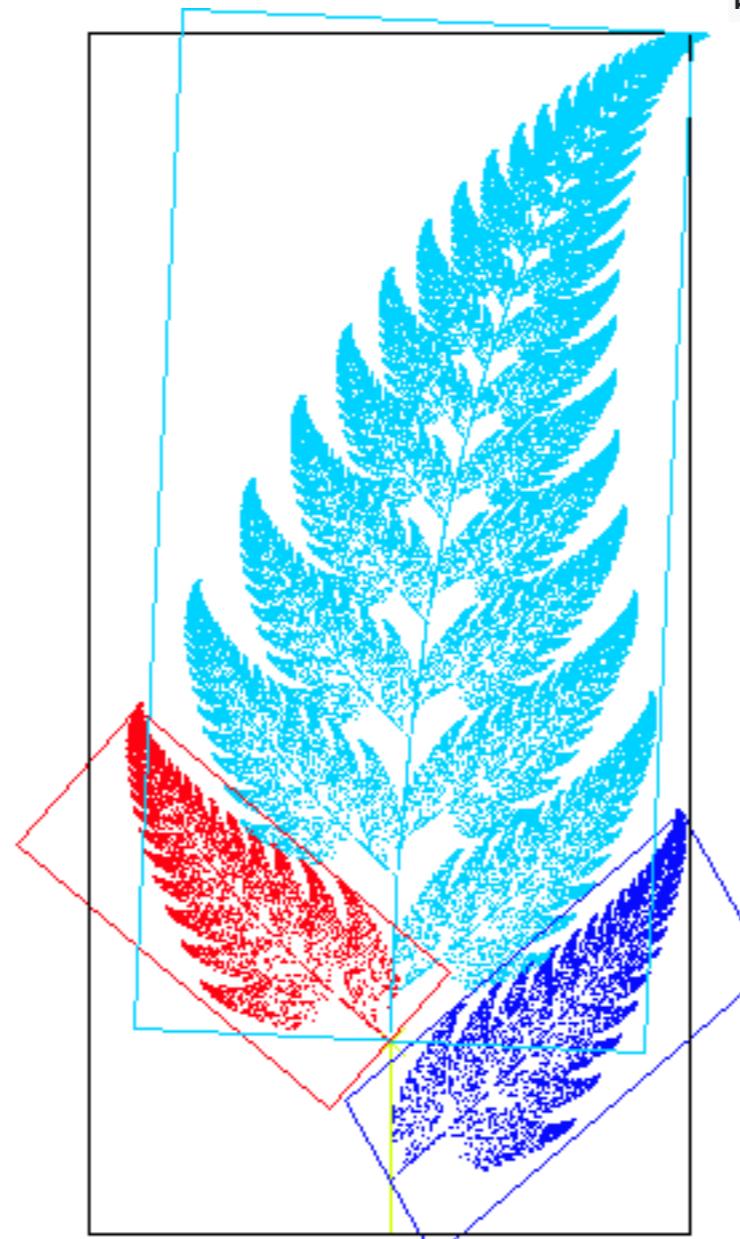
Self-similarity

fractals

Self-organization

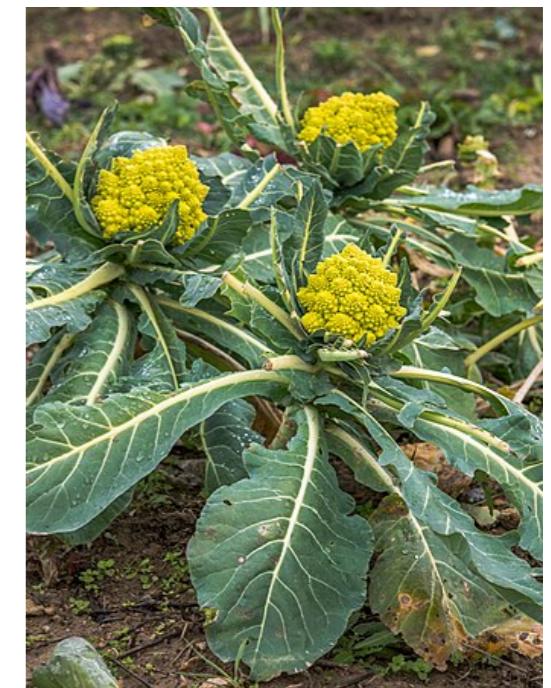


Hausdorff dimensionality
Poincare section
Stroboscopic sampling



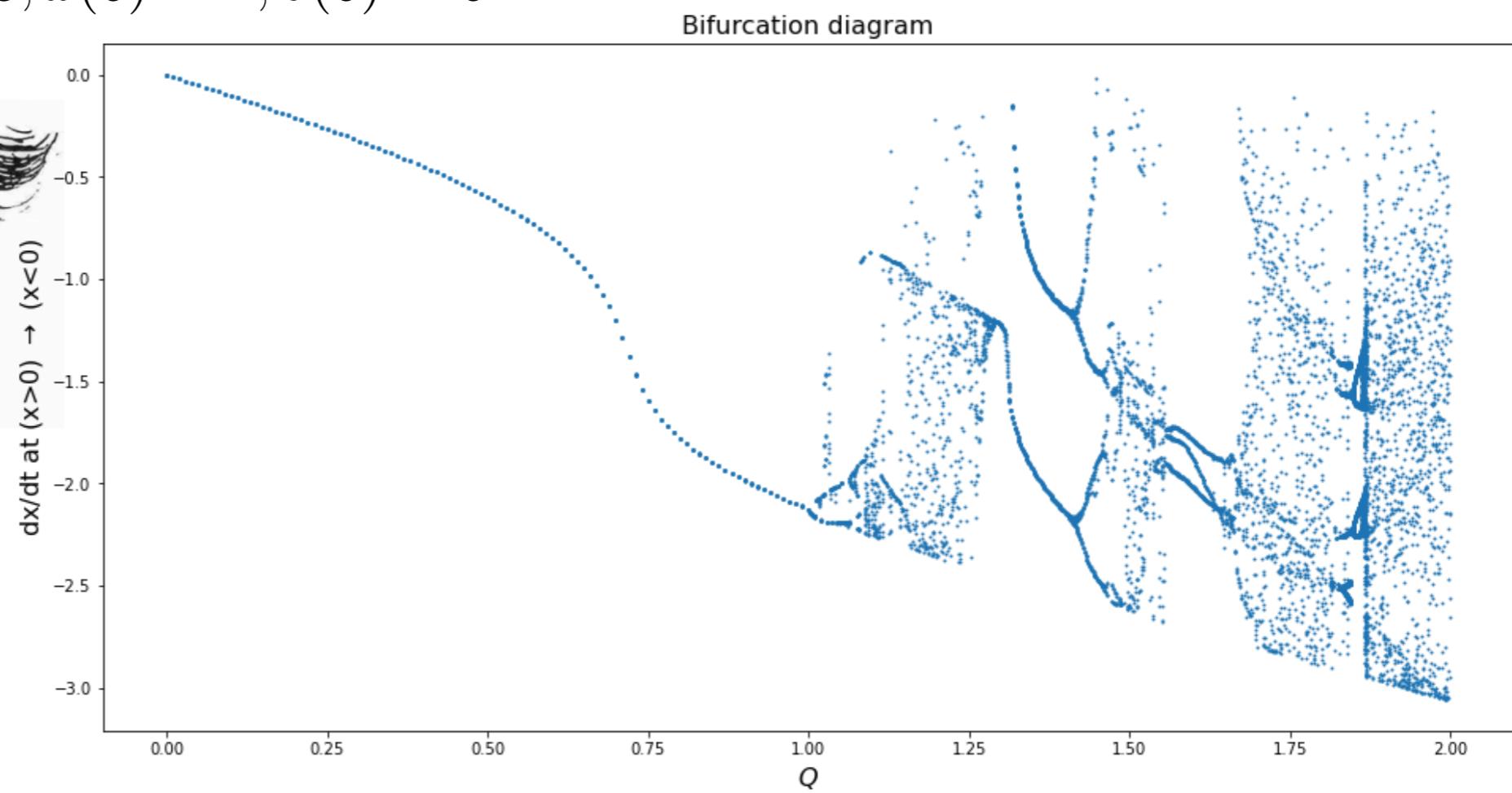
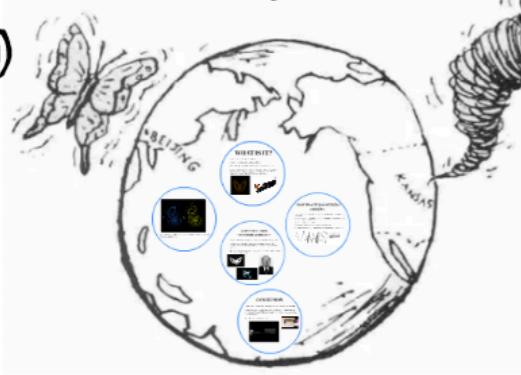
Romanesco broccoli

西兰花



$$k = 1, \gamma = 0.5, \Omega = 2/3, x(0) = 1, v(0) = 0$$

BUTTERFLY EFFECT (CHAOS THEORY)



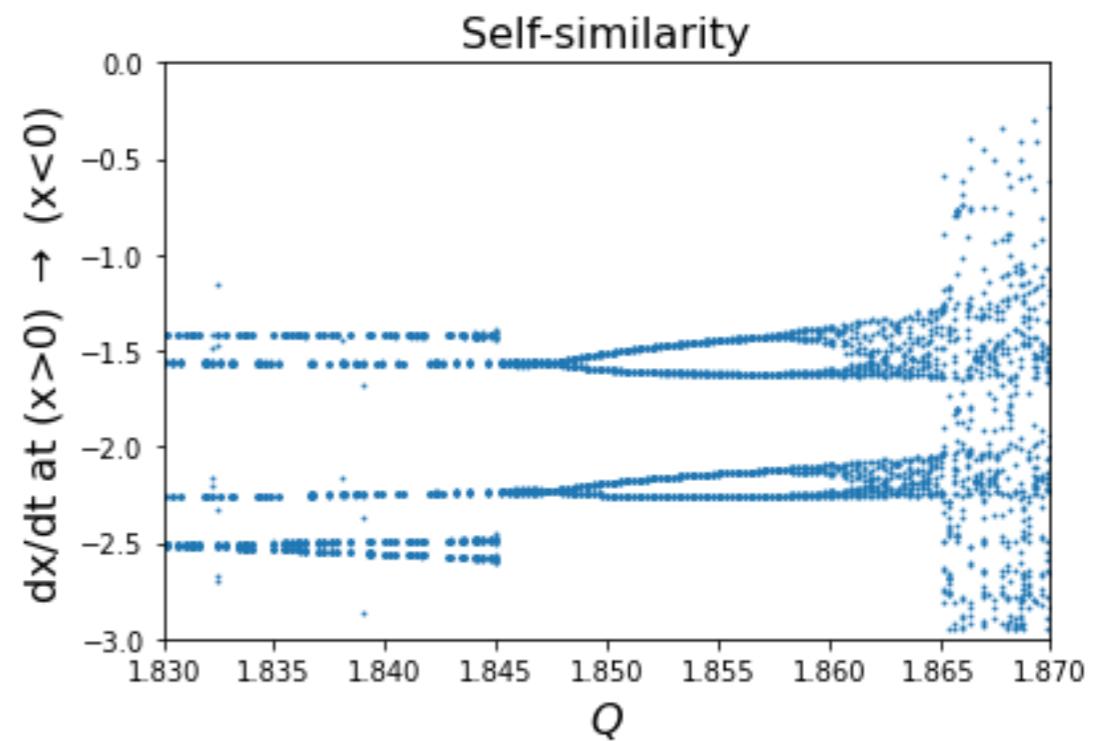
Record the velocity as x passes 0 from above

Periodic motion, discrete set of period

$$nT_\Omega$$

Chaotic motion, many distinct values

**Hausdorff dimensionality
Poincare section
Stroboscopic sampling**



Many interesting videos

<https://www.myphysicslab.com/pendulum/chaotic-pendulum-en.html>

Spectral Analysis and Power Spectrum

$$x(t) \quad t_i = i \cdot \tau : i = 0, 1, \dots, N-1, \quad \tau = \frac{T}{N-1}$$

$$g(\omega) \quad \omega_k = 2\pi k/(N\tau) : k = 0, 1, \dots, N-1, \quad \Delta\omega = 2\pi/(N\tau)$$

$$g(\omega_k) = \sum_{i=0}^{N-1} e^{-i\omega_k \cdot t_i} x_i = \sum_{i=0}^{N-1} e^{-i\frac{2\pi k}{N\tau} i\tau} x_i = \sum_{i=0}^{N-1} (W_N)^{ik} x_i \quad W_N = e^{-i2\pi/N}$$

$$g_k = \sum_{i=0}^{N-1} (W_N)^{ik} x_i, \quad k = 0, 1, \dots, N-1$$

Starting from N time points, we have N frequencies

$$x_i = \frac{1}{N} \sum_{k=0}^{N-1} (W_N^*)^{ki} g_k, \quad i = 0, 1, \dots, N-1$$

$$= \frac{1}{N} \sum_{k=0}^{N-1} \sum_{j=0}^{N-1} (W_N^*)^{ki} (W_N)^{jk} x_j = \sum_{j=0}^N \left(\frac{1}{N} \sum_{k=0}^{N-1} e^{-i2\pi k(i-j)/N} \right) x_j = \sum_{j=0}^{N-1} \delta_{i,j} x_j = x_i$$

Direct Fourier Transformation: computational complexity $O(N^2)$

Fast Fourier Transformation: computational complexity $O(N \log_2 N)$



Chap.12.

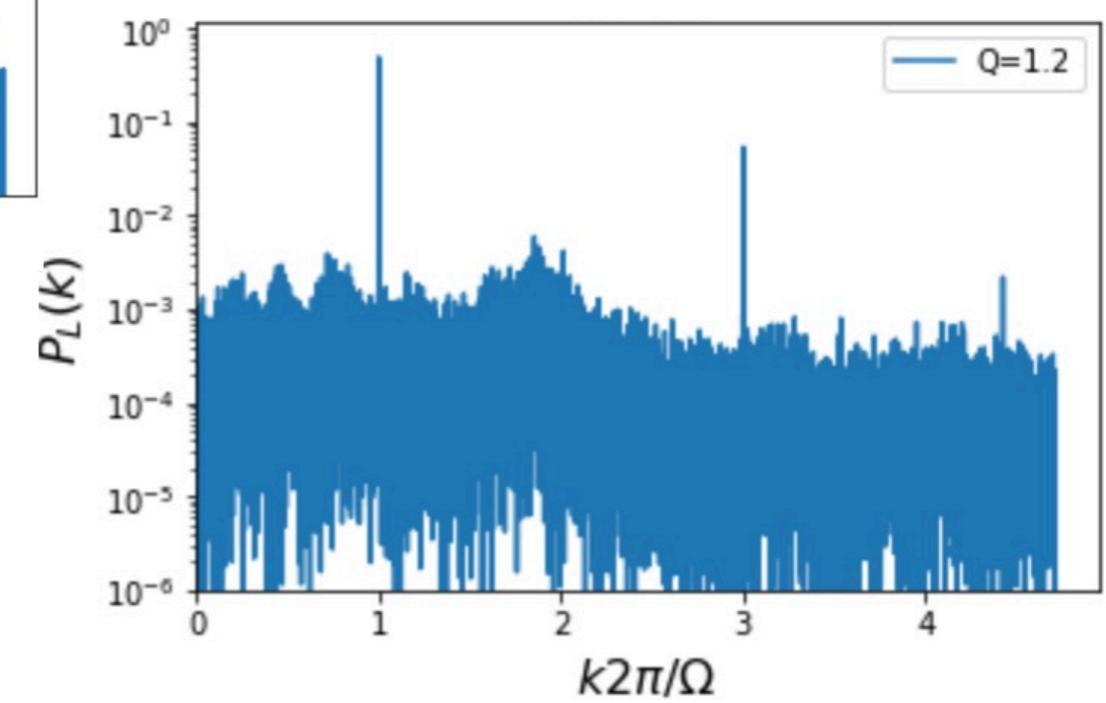
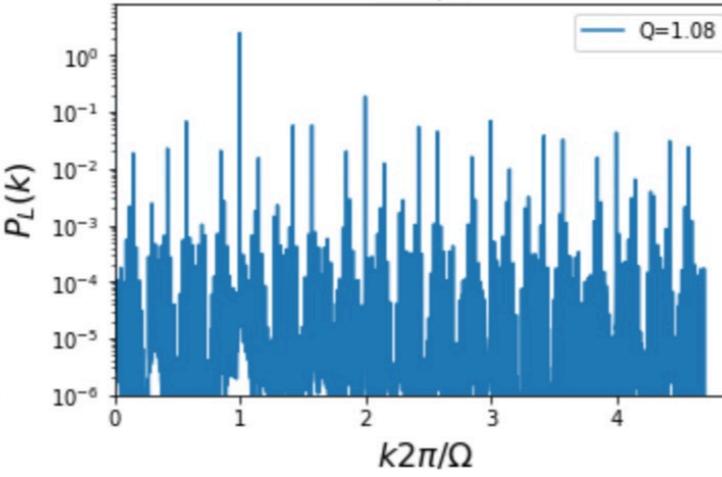
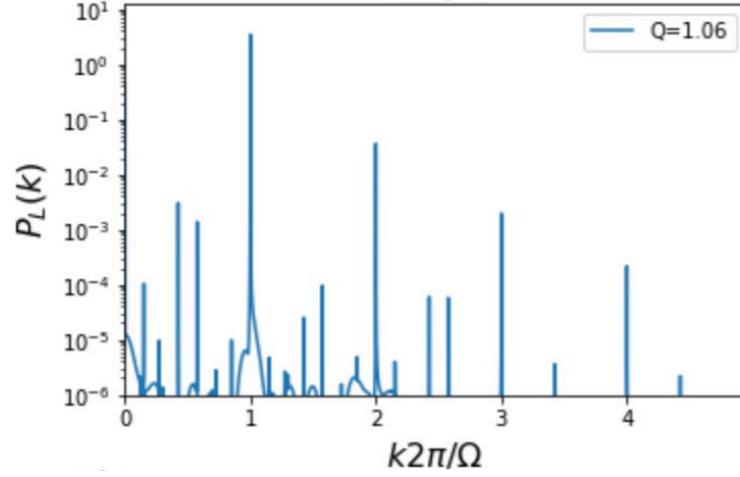
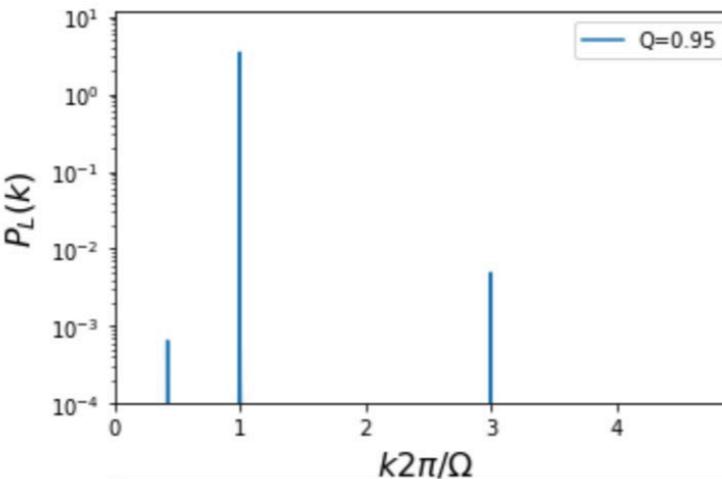
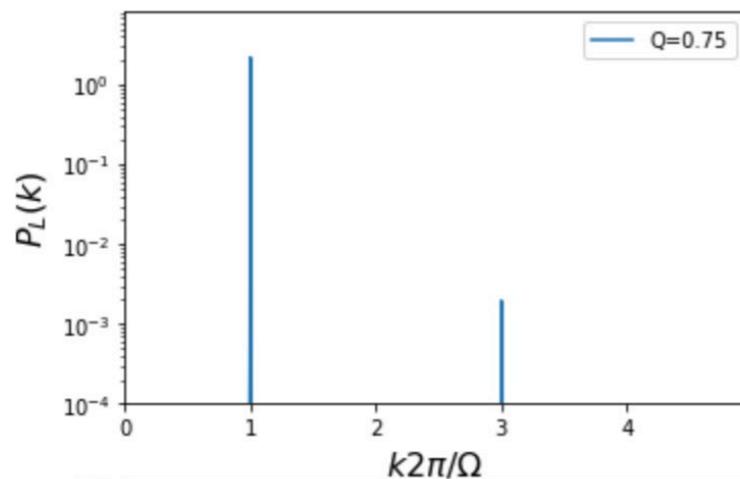
Power spectrum

space \longleftrightarrow momentum
time \longleftrightarrow frequency

$$P(0) = \frac{1}{N^2} |g_0|^2$$

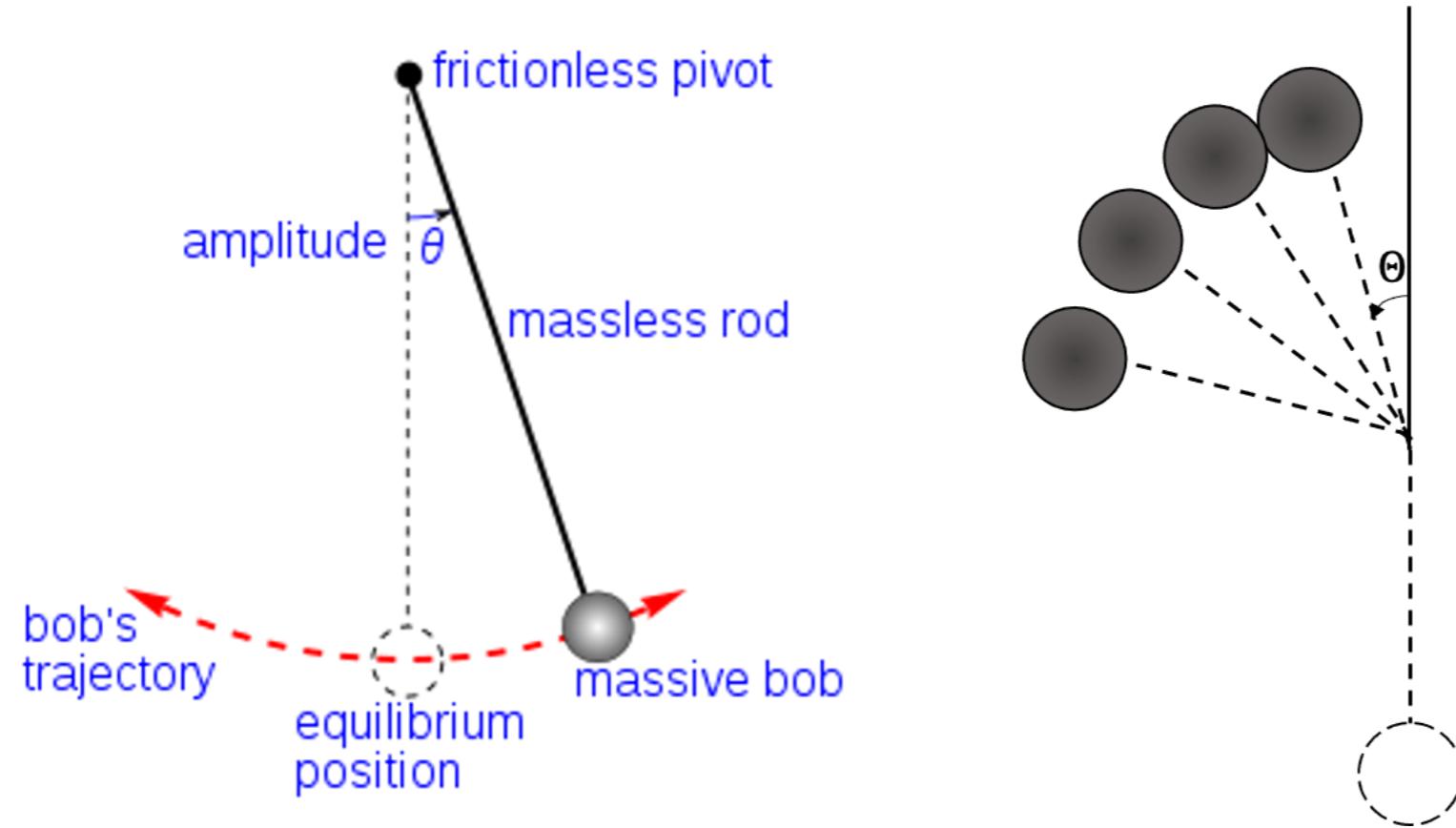
$$P(k) = \frac{1}{N^2} [|g_k|^2 + |g_{N-k}|^2], \quad k = 1, 2, \dots, \frac{N}{2} - 1$$

$$P\left(\frac{N}{2}\right) = \frac{1}{N^2} |g_{\frac{N}{2}}|^2 \quad \text{assume } N \text{ is even}$$



$$T = \frac{2\pi}{\Omega}$$

Soliton in pendulum

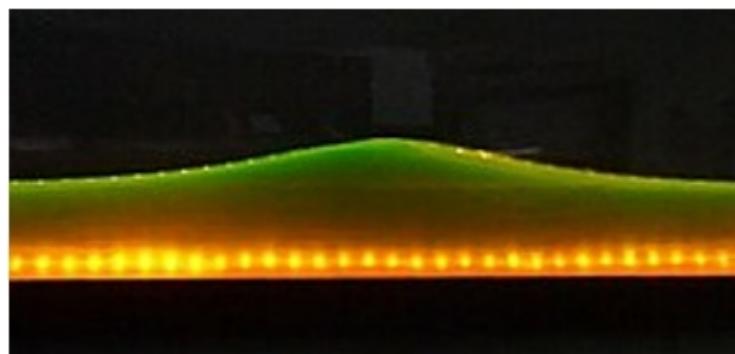


$$\theta = \theta + \pi$$

$$\ddot{\theta} = \frac{g}{l} \sin(\theta)$$

The ball will spend most of its time on the peak while rapidly go through the other region. Showing a well localized excitation at temporal space — Soliton or Instanton.

Classics soliton : water wave



BPST instanton: solution of equation of motion of SU(2) Yang-Mills theory in Euclidean space-time with winding-number 1. Meaning it describe the transition between two topologically different vaccum.

